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Abstract: We present results relating to the existence and uniqueness of the free energy and equilibrium states for classical continuous spin systems with superstable interactions.

We consider the lattice  $\mathbb{Z}'$  at each site of which there is a vector spin variable S,  $x \in \mathbb{Z}'$ ,  $S \in \mathbb{R}'$ . We denote by  $\underline{S} \in \{S\}$  a spin configuration on  $\mathbb{Z}'$ . Each S has associated with it an intrinsic positive measure  $\mu(dS)$ , the same for all sites, such that  $J\mu(dS)$  e<sup> $-\alpha S^{2}$ </sup> <  $\infty$  for  $\alpha > 0$ . The energy of a given spin configuration  $S_{\Lambda}$  in  $\Lambda \subseteq \mathbb{Z}'$  consists of both pair and self interactions and satisfies the following conditions:

a) Superstability There exists A > 0,  $C \in R$  such that

$$U(S_{\Lambda}) \geq \sum_{\mathbf{x} \in \Lambda} [A S_{\mathbf{x}}^{2} - C]$$
 (1)

where  $S_{\Lambda}$  is a configuration in  $\Lambda_{ullet}$ 

b) Regularity If  $^{\Lambda}_{1}$ ,  $^{\Lambda}_{2}$  are disjoint then their interaction energy  $W(S_{\Lambda_{1}} | S_{\Lambda_{2}}) = U(S_{\Lambda_{1}} | S_{\Lambda_{2}}) - U(S_{\Lambda_{1}}) - U(S_{\Lambda_{2}})$  has the bound

$$|W(S_{\Lambda_{1}}|S_{\Lambda_{2}})| \leq \frac{1}{2} K \sum_{\mathbf{x} \in \Lambda_{1}} \Sigma |S_{\mathbf{x}}| |S_{\mathbf{y}}| |\mathbf{x}-\mathbf{y}|^{-\nu-\varepsilon}$$
(2)

where  $|\mathbf{x}| = \max_{1 \le i \le v} |\mathbf{x}^i|, |\mathbf{S}| = \begin{bmatrix} \frac{d}{\Sigma} & (\mathbf{S}^i)^2 \end{bmatrix}^{\frac{1}{2}}.$ 

For  $\Lambda$  bounded in  $\mathbb{Z}'$  we consider the restriction, S of S to  $\Lambda^c$  and define the partition function  $Z(\Lambda|S)$  and free energy per site  $F(\Lambda|S)$  with boundary conditions' (b.c.) determined by S as

$$Z(\Lambda | S_{\Lambda^c}) = \int \mu_{\Lambda}(dS_{\Lambda}) \exp \left[-U(S_{\Lambda}) - W(S_{\Lambda} | S_{\Lambda^c})\right]$$

$$F(\Lambda | S_{\Lambda^c}) = |\Lambda|^{-1} l_n Z(\Lambda | S_{\Lambda^c})$$

where  $\mu_{\Lambda}(dS_{\Lambda}) = \prod_{\mathbf{x} \in \Lambda} \mu(dS_{\mathbf{x}})$ ,  $|\Lambda| = \#$  of sites in  $\Lambda$ . (The dependence on temperature and magnetic field is included in U and W; it will be made explicit when necessary.)

Theorem 1. Let (1) and (2) hold and let  $\underline{S} \in \mathcal{H}_a$ :

Remark: "Zero" b.c. correspond to S=0 for  $x \in \Lambda^{C}$ . The thermodynamic limit of the "periodic" b.c. free energy can also be shown to exist and be equal to F.

A probability measure  $\vee$  on the configuration space  $\{S\}$  is said to be regular if it satisfies the following condition: There exists  $\gamma \geq 0$ ,  $\delta \geq 0$ , such that for every  $\Delta$  bounded in  $Z^{\nu}$  and  $N^2 \geq 0$  the following holds:

$$v[B(N^2|\Delta)] < \exp[-|\Delta| (\gamma N^2 - \delta]$$

where

$$B(N^{2}|\Delta) = \{\underline{s} \mid \sum_{x \in \Delta}^{:} s_{x}^{2} \geq N^{2} |\Delta| \}$$

For  $\Lambda$  bounded in  $Z^{\vee}$  we denote by  $F(\Lambda|\nu)$  the free energy in  $\Lambda$  for boundary conditions specified by the measure  $\nu$  as

$$F(\Lambda | \nu) = \int \nu(d\underline{s}) F(\Lambda | \underline{s}_{\Lambda^c})$$

Theorem 2. Let (1) and (2) hold and let  $\Lambda$  be as in theorem 1 then for  $\nu$  regular  $\lim_{n\to\infty} F(\Lambda_n|\nu) = F$ .

The finite volume equilibrium measure with boundary conditions S  $_{\Lambda^c}$ ,  $_{\Lambda^c}$  is given by

$$v_{\Lambda}(\mathrm{d}s_{\Lambda}|s_{\Lambda^{\mathbf{c}}}) = z^{-1} (\Lambda|s_{\Lambda^{\mathbf{c}}}) \mu(\mathrm{d}s_{\Lambda}) \exp[-u(s_{\Lambda}) - w(s_{\Lambda}|s_{\Lambda^{\mathbf{c}}})]$$
 (3)

a measure  $\nu$  on  $\{S\}$  is said to be an equilibrium measure (for our system) if its conditional probabilities  $\nu(dS_{\Lambda}|S)$  satisfy the Dobrushin, Lanford and Ruelle (DLR) [2] equations, i.e. eq. (3).

Theorem 3. Let the conditions of Theorem 1 be satisfied and let  $\bigvee_{n} (dS_{n} | S_{n})$ 

be finite volume equilibrium states then it is always possible to choose subsequences  $\{n_i\}$  (which may depend on the b.c.) such that  $\bigvee_{\substack{\Lambda \\ n_i}} (ds_{\Lambda} \mid S) \rightarrow \bigvee_{\substack{\Lambda \\ n_i}} (ds_$ 

a <u>regular</u> equilibrium measure on {S}, [3].

The one component spin system, S  $\in$  R, will be called ferromagnetic with translation invariant interactions if

$$U(S_{\Lambda}) = -\frac{1}{2} \sum_{x \neq y \in \Lambda} J(x-y) S_{x}S_{y} - \sum_{x \in \Lambda} \Phi(S_{x}) - h \sum_{x \in \Lambda} S_{x}; J(x) \geq 0.$$

Theorem 4. Let v be a regular equilibrium measure of a ferromagnetic system in an external field h whose interactions satisfy (1) and (2) then v is unique (and hence translation invariant) whenever the infinite volume free energy F(h)

is differentiable with respect to h [4].

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## References

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