

Weakly interacting Bose gas in disordered environment*

G.M. Falco and T. Nattermann
University of Cologne, Germany

V.L. Pokrovsky,
TAMU, TX, USA & Landau Institute , Moscow

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Outline

Ideal Bose gas in random potential

Weakly repulsive gas in a random potential

Bosons in one dimension

Bosons in traps

Introduction

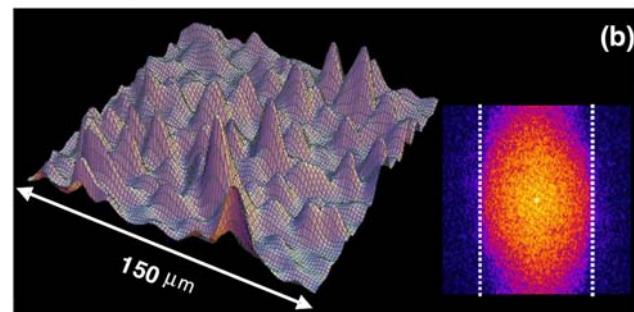
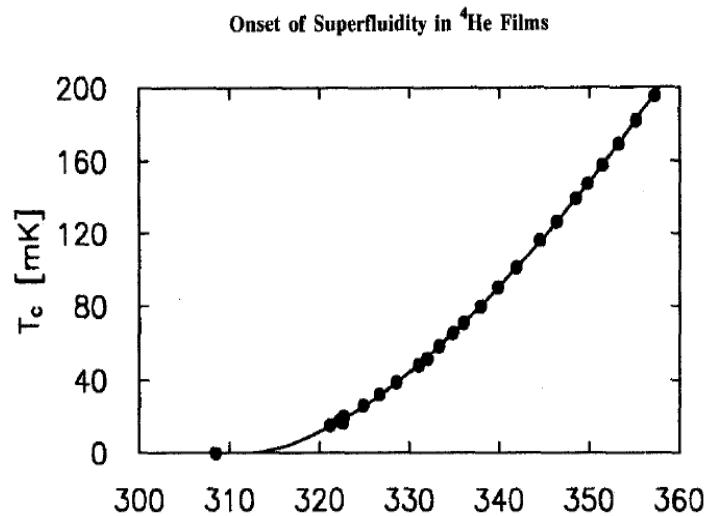
BEC : finite part of atoms in the state with minimal energy.

Examples: Superfluid ^4He , laser cooled atoms in a trap

Disorder: Superfluid He in porous media (J.D. Reppy et al '92)
Cold atoms in speckle potential (R.G. Hulet et al. '08)

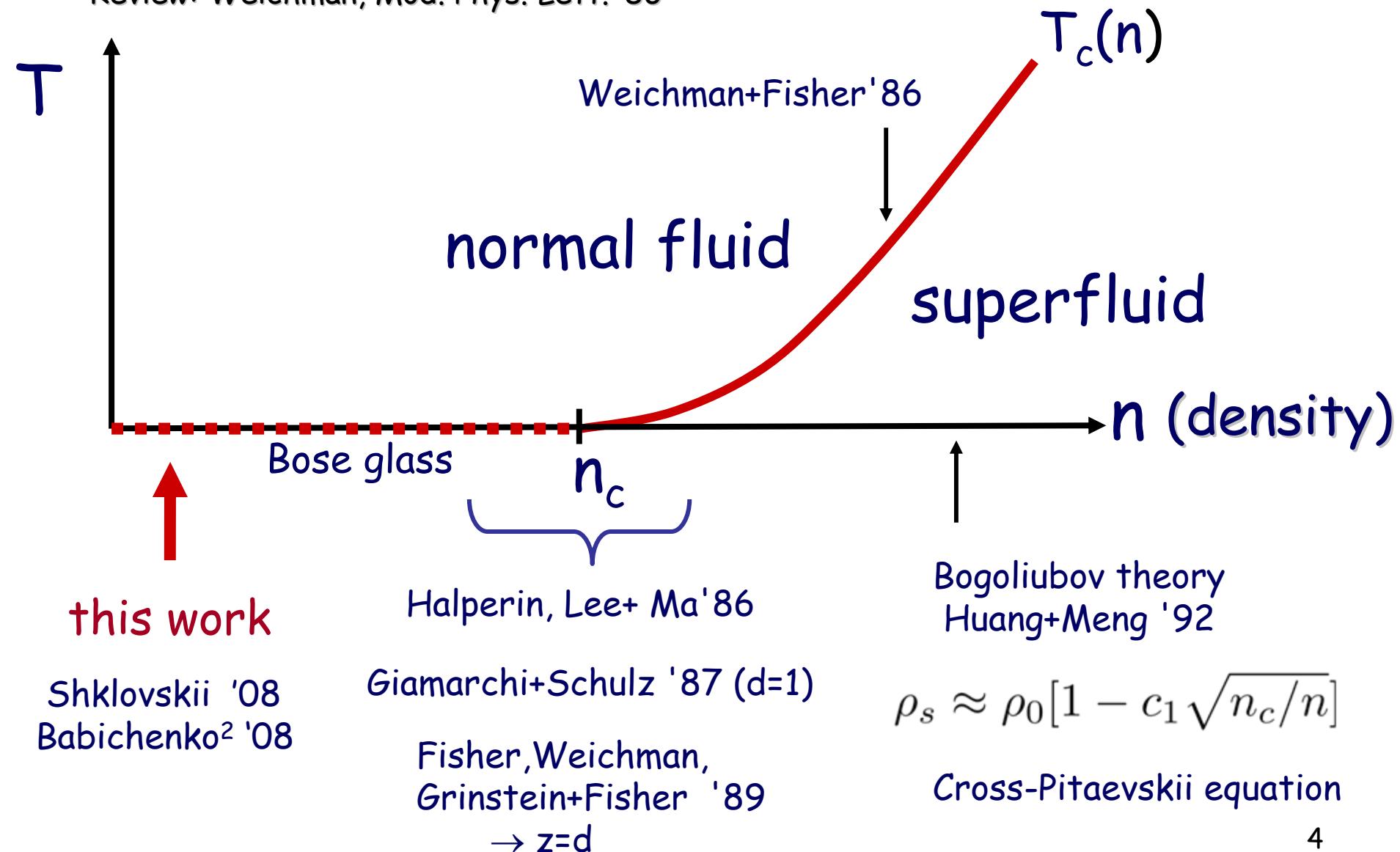


Breakdown of superfluidity at strong disorder



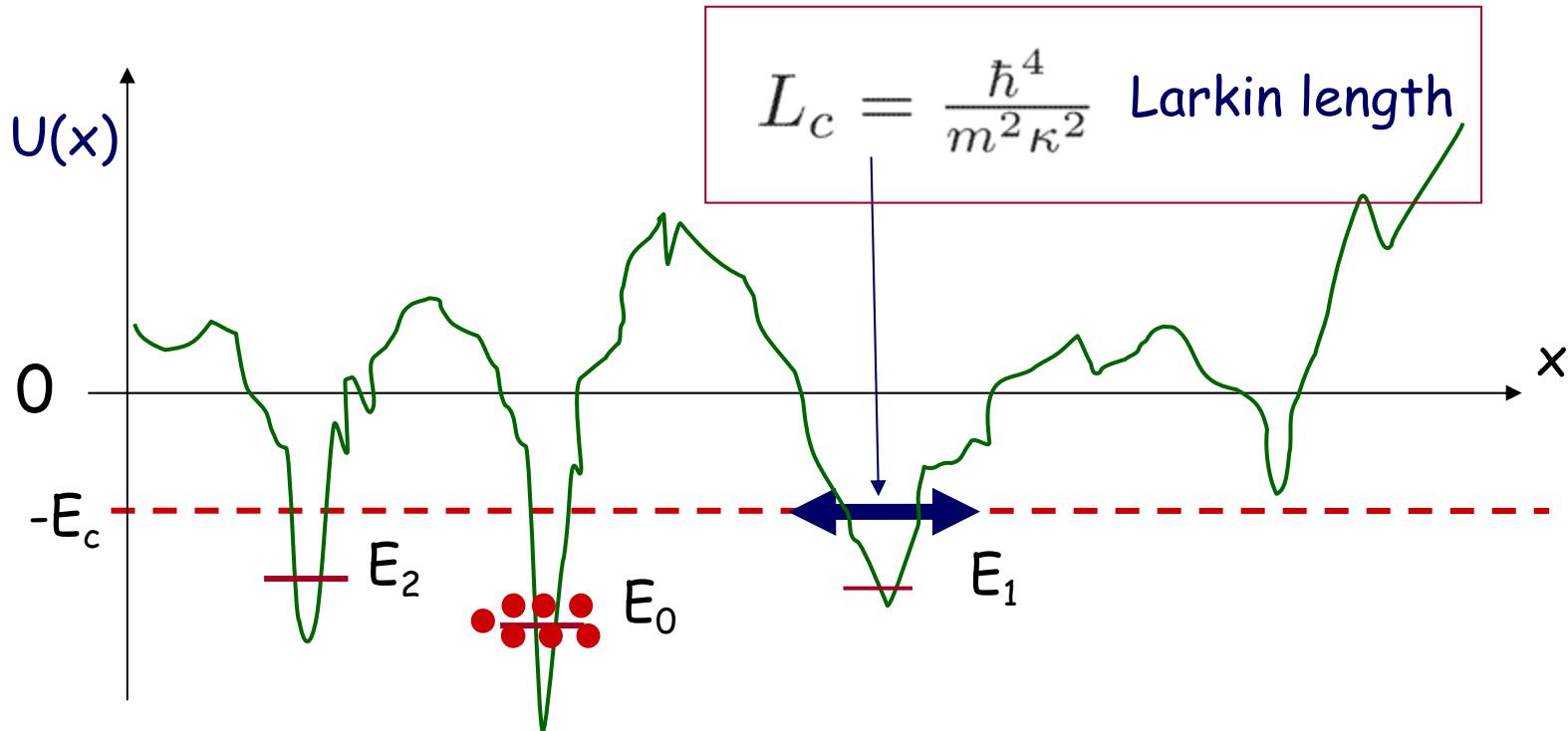
Bosons in disordered environment:

Review: Weichman, Mod. Phys. Lett. '86



Ideal 3d Bose gas in random potential

$$\frac{\hbar^2}{2m} \nabla^2 \psi + (E - U(\mathbf{x})) \psi = 0 \quad \langle U(\mathbf{x})U(\mathbf{x}') \rangle = \kappa^2 \delta(\mathbf{x} - \mathbf{x}')$$



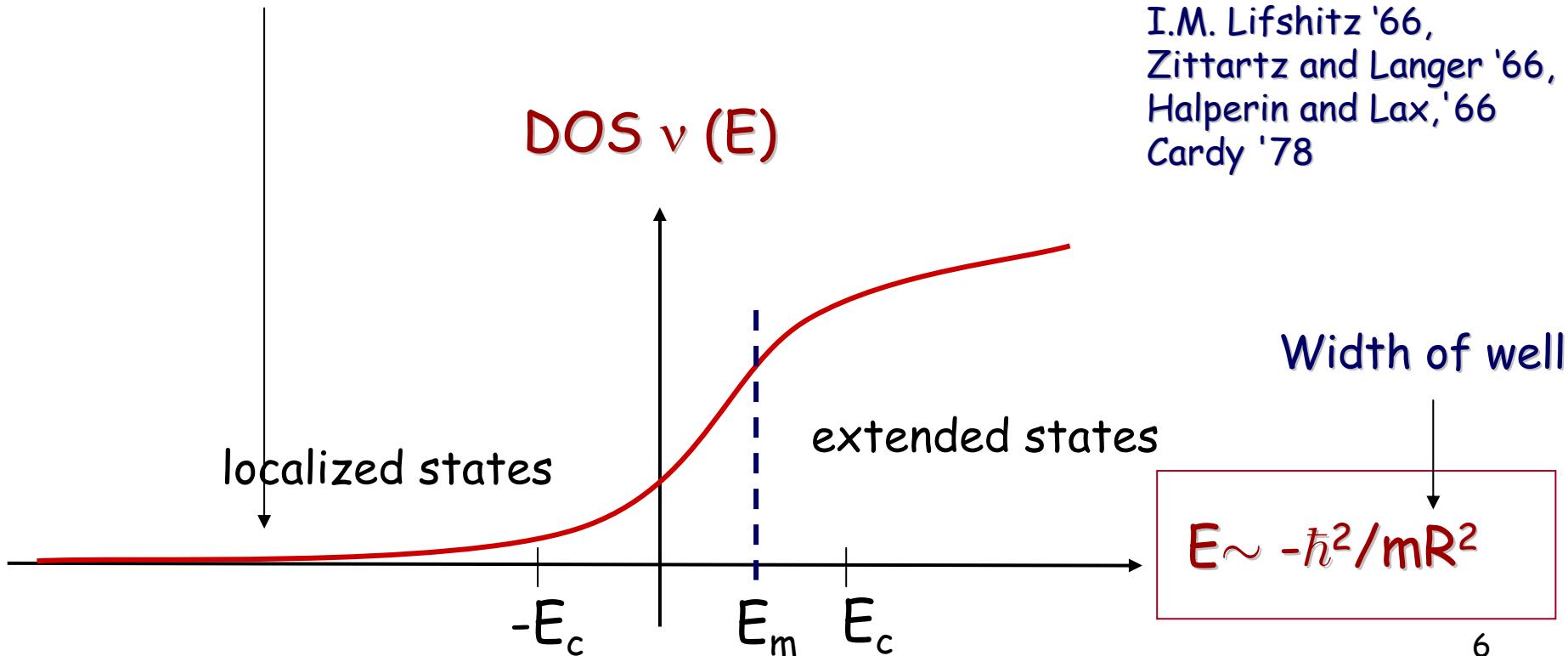
$$E_c = \frac{\hbar^2}{2m L_c^2} \quad \text{Larkin energy}$$

T=0: All particles in ground state $E_0 \approx -E_c \ln^2(L_0/L_c)$

Ideal Bose gas in random potential

DOS for $E \ll -E_c$ dominated by wells of width $R \sim \hbar/\sqrt{m|E|} \ll L_c$

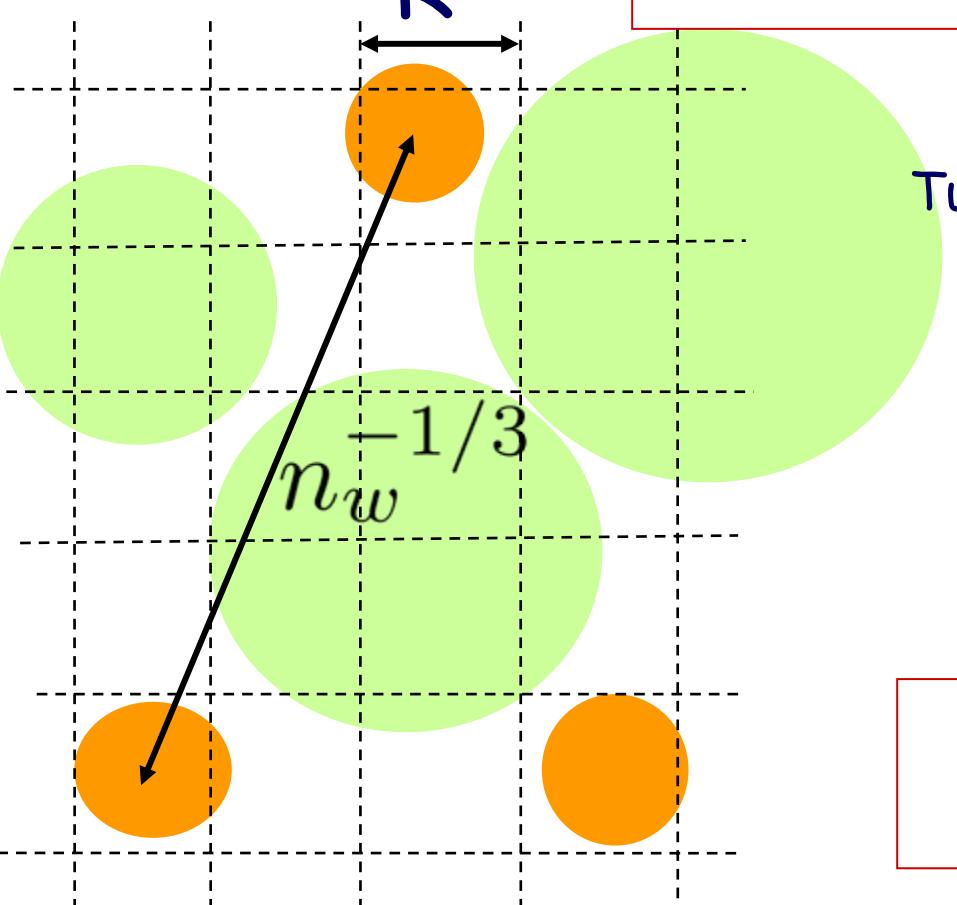
$$\nu(E) = \frac{1}{V} \langle \delta(E - E[U(\mathbf{x})]) \rangle \sim |E|^{3/2} e^{-\sqrt{|E|/E_c}}$$



Ideal Bose gas in random potential

Spatial density $n_w(R)$ of wells with radius $\ll R \ll L_c$ ($E \ll -\hbar^2/(2mR^2) \ll E_c$)

$$n_w(R) = \int_{-\infty}^{-\frac{\hbar^2}{2mR^2}} dE \nu(E) \sim \frac{L_c}{R^4} e^{-L_c/R}$$



Tunneling amplitude $t(R)$ between wells with radius $\ll R$:

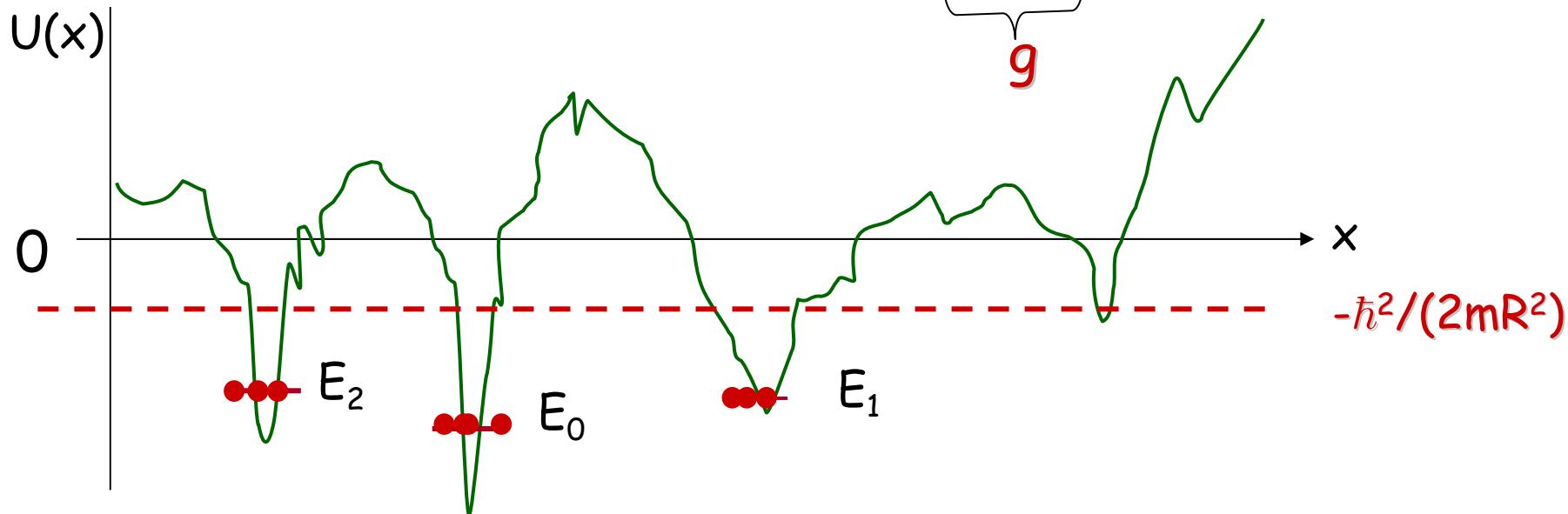
$$t(R) = \exp\left(-\frac{1}{\hbar} \int |p| dl\right)$$

$$\frac{1}{\hbar} \int |p| dl \approx n_w^{-1/3}/R \sim e^{L_c/3R}$$

$$t(R) \sim e^{-\left(\frac{R}{L_c} e^{L_c/R}\right)^{1/3}}$$

Weakly repulsive bosons in a random potential

$$\mathcal{H} = \int d^3x \Psi^\dagger \left(-\frac{\hbar^2}{2m} \nabla^2 + U(\mathbf{x}) + \frac{2\pi\hbar^2 a}{m} \Psi^\dagger \Psi \right) \Psi$$



Assume that all potential wells with radii up to R are filled:

\Rightarrow number of particles per well of size R : $N_w(R) = n/n_w(R) \gg 1$

\Rightarrow repulsion energy per particle: $E_g(R) \approx g N_w/R^3 \sim g n e^{L_c/R}$

\Rightarrow total energy per particle: $\mu(R) = -\hbar^2/(2mR^2) + E_g(R)$

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Minimization over R : $\Rightarrow R(n) = L_c / \ln(n_c/n)$, $n \ll n_c \approx 1/(3L_c^2a)$

$$\mu(n) = -\frac{\hbar^2}{2mR^2(n)} = -\frac{1}{2}E_c \left(\ln \frac{n_c}{n} \right)^2$$

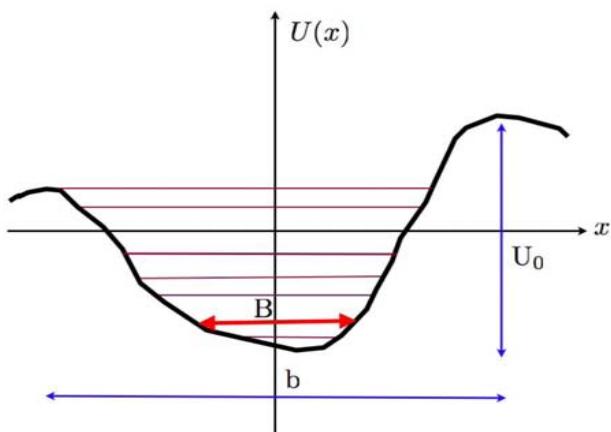
$$\frac{n_c}{n} = \frac{\xi^2}{L_c^2} \quad (\text{Fermions: } L_c \rightarrow a)$$

Preliminary conclusions

- ⇒ At $n \ll n_c$ Bose gas decays into fragments,
particle density in fragments each of density $n_c \sim 1/(aL_c^2)$
- ⇒ tunneling exponentially suppressed: $t(n) \sim e^{-c(n_c/n)^{1/3}}$
- ⇒ particle number in fragments $N_w = L_c / [3a(\ln \frac{n_c}{n})^3]$ well defined
- ⇒ phase uncertain, no phase coherence ⇒ no superfluidity
- ⇒ finite compressibility $\frac{n}{E_c} \ln \left(\frac{n_c}{n} \right)$ „Bose glass“
- ⇒ $\hat{H}_{\text{eff}} = \sum_j C_j (\hat{N}_j - \langle N_j \rangle)^2 - \sum_{i,j} t_{ij} \cos(\hat{\phi}_i - \hat{\phi}_j)$
- ⇒ charged bosons VRH $\sigma(T) \sim e^{-C[E_c n_c / (Tn)]^{1/4}}$

For $n \approx n_c$ i.e. fragments merge → transition to superfluid

Correlated disorder



$$\langle U(\mathbf{x})U(\mathbf{x}') \rangle = \frac{U_0^2}{b^3} e^{-|\mathbf{x}-\mathbf{x}'|/b}$$

\Rightarrow 2 length scales $b, B \sim \hbar^2/mU_0)^{1/2}$

$b \ll B \Rightarrow$ uncorrelated disorder

$b \gg B \Rightarrow$ new results

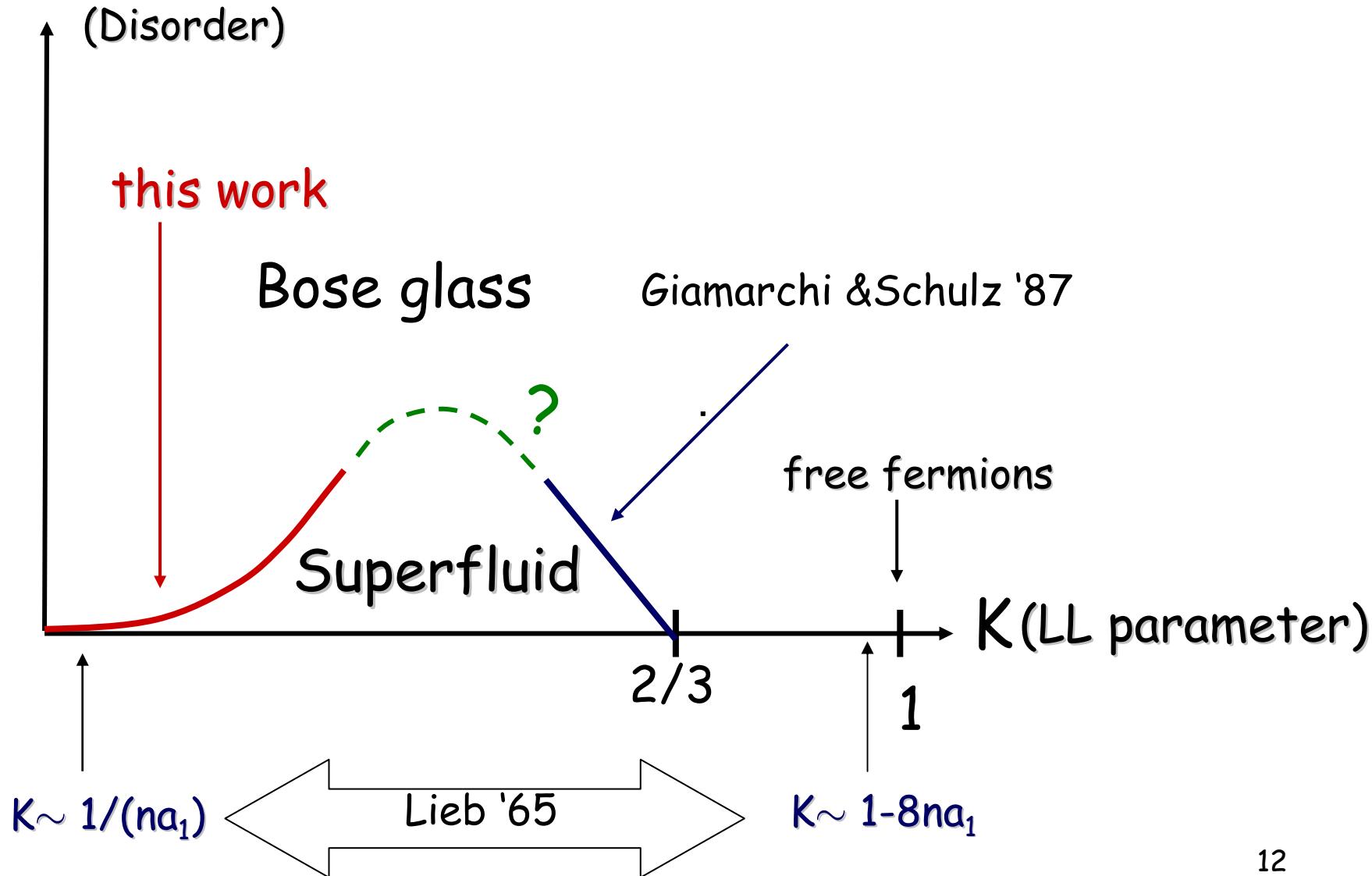
$$\nu(E) \sim |E|^3 \exp(-E^2/2U_0^2)$$

Keldysh & Proshko '63
Kane '63
Shklovskii and Efros '70
John & Stephen '84

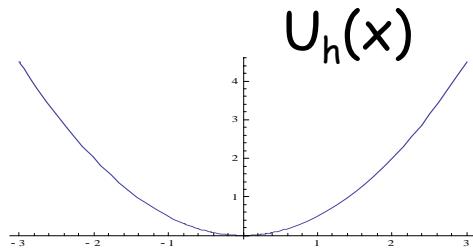
$$\mu(b, n) \approx -U_0 \sqrt{2 \ln(\frac{n_c}{n})}$$

$$n \ll n_c \sim 1/(B^2 a)$$

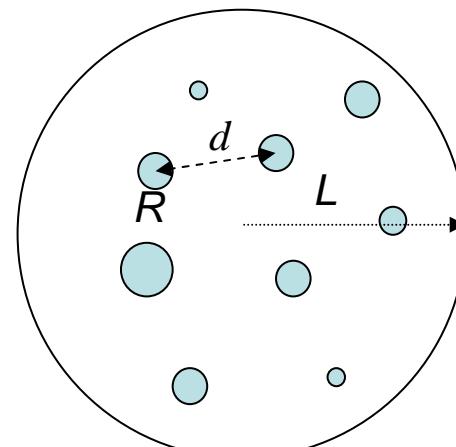
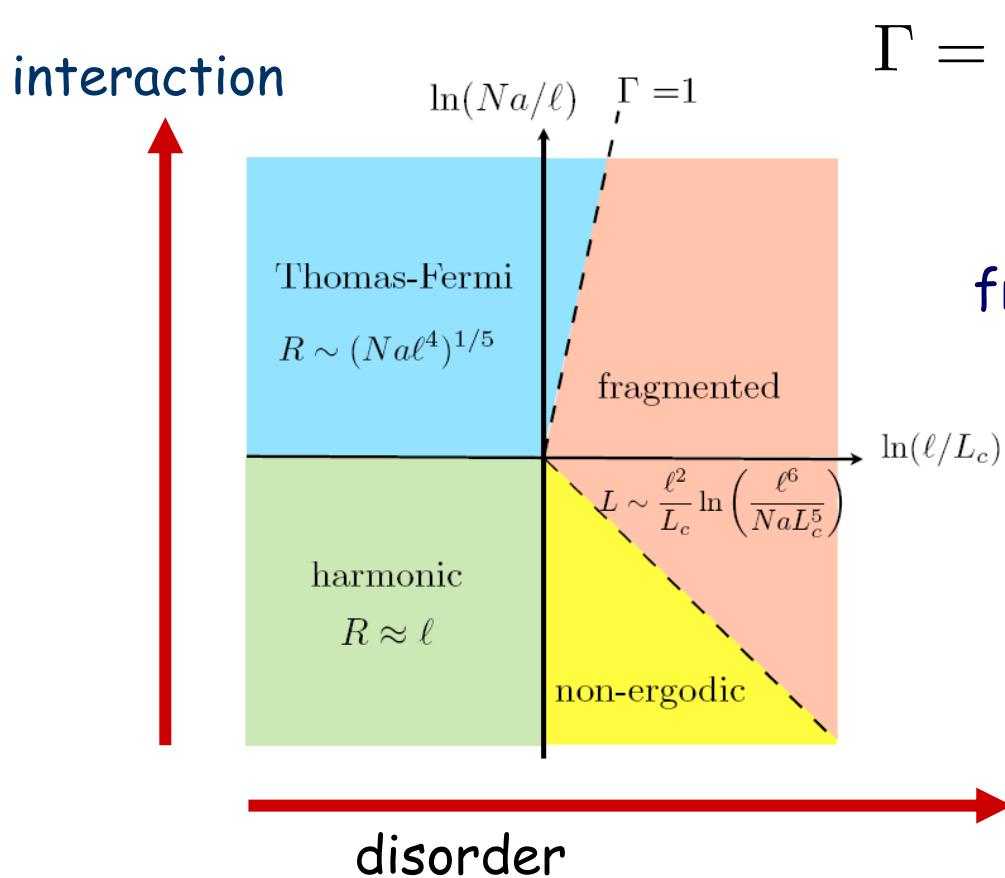
Bose gas in one dimensions



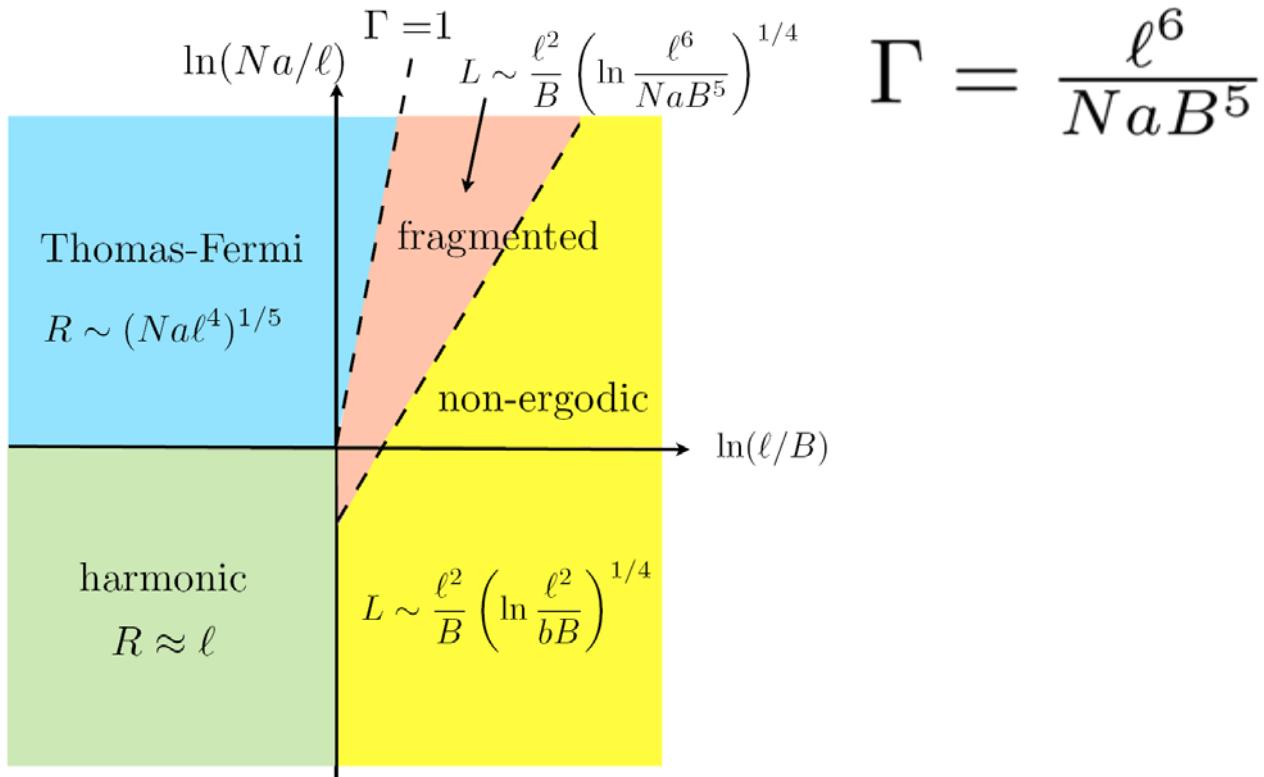
Bosons in traps (uncorrelated disorder)



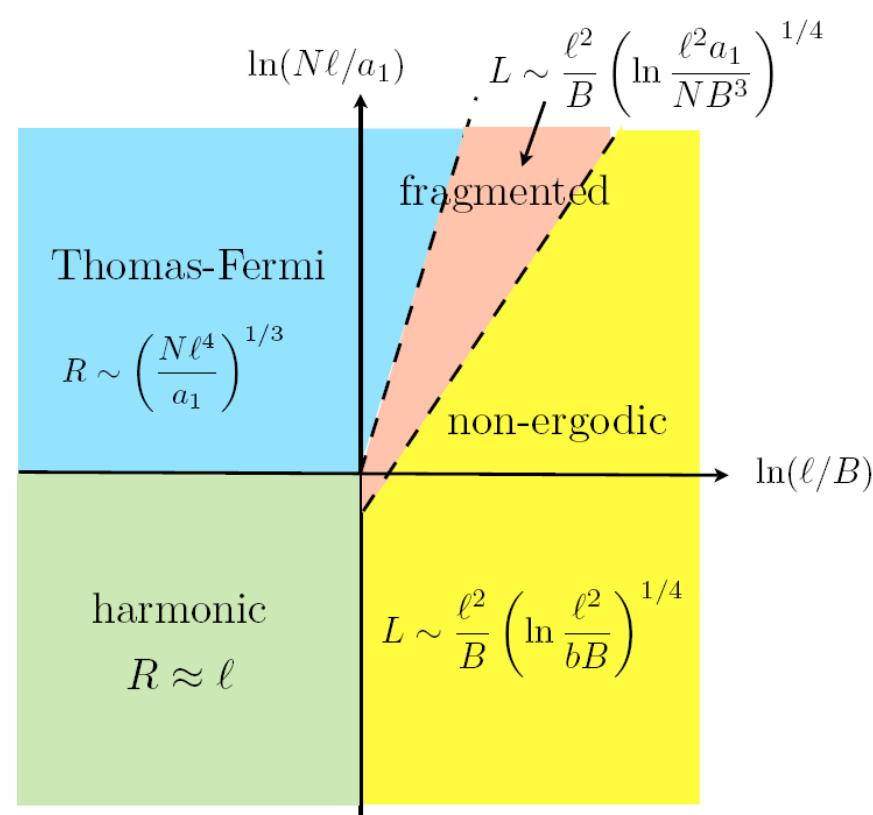
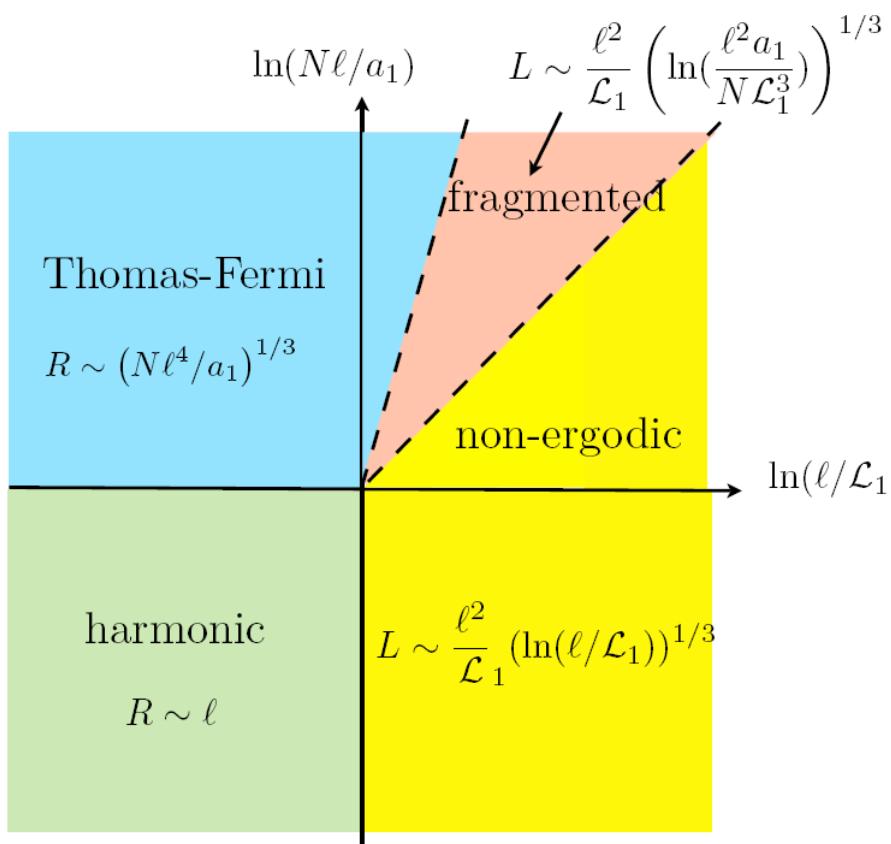
$$\mu(R) = -\frac{\hbar^2}{2mR^2} + E_{int}(R) + \frac{\hbar^2}{2m} \frac{R^2}{\ell^4}$$



Bosons in traps (correlated disorder, d=3)



Bose gas in 1 dimensions: parabolic trap



Conclusions

- Semi-quantitative analysis of the phase states of a weakly interacting strongly diluted Bose gas in a random Gaussian potential.
- The system is characterized by the Larkin length L_c and the scattering length a
- At particle density $n \ll n_c \approx 1/(aL_c^2)$ the Bose particles occupy deep potential wells and exponentially weakly tunnel to other wells. The number of particles in each well is defined, but phases are uncertain.
- At average particle density $n \approx n_c$ the transition to the superfluid proceeds.
- In a trap the oscillator length l appears as a new length scale. Four different regimes are found, depending on the mutual strength of L_c , aN and l , respectively.
- All results can be extended to lower dimensions and to correlated disorder.

Theory

Corrections to the Bogolyubov theory caused by a weak disorder: $n \gg n_c$

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Possible Bose-glass state

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R.T. Scalettar et al., Phys. Rev. Lett. **66**, 3144 (1991)
W. Krauth et al., Phys. Rev. Lett. **67**, 2307 (1991)

Spin model, transition from the normal state to superfluid

M. Ma, B.I. Halperin and P.A. Lee, Phys. Rev. B **34**, 3136 (1986)

One-dimensional Hubbard model

T. Giamarchi and H. Schulz, Europhys. Lett. **3**, 1287 (1987) $n \gg n_c$